

Today's Talk

This talk will introduce recent development in an asymptotic expansion approach in finance, particularly, the following topics:

- **asymptotic expansion in a general diffusion setting for approximations of density functions and option prices**
- **asymptotic expansion for basket option pricing in a local-stochastic volatility (LSV) with jumps**
- **perturbative expansion method for backward stochastic differential equations (BSDEs)**
- **perturbation technique with interacting particle method for BSDEs**
- **(On an application to of the method to mean-variance hedging problems in partially observable markets with stochastic filtering, please see “Making Mean-Variance Hedging Implementable in a Partially Observable Market,” *Quantitative Finance*, published online: 20 Mar 2014.)**



- Let \mathcal{S} be the real Schwartz space of rapidly decreasing C^∞ -functions on \mathbf{R} and \mathcal{S}' be its dual space.
- Next, take $\Phi \in \mathcal{S}'$. Then, the asymptotic expansion of $\Phi(G^{(\epsilon)})$ as $\epsilon \downarrow 0$ can be verified by Watanabe theory. (e.g. Watanabe(1987))
- In particular, if we take the delta function at $x \in \mathbf{R}$, δ_x as Φ , we obtain an asymptotic expansion of the density for $G^{(\epsilon)}$.
- We define the following notations:

$$\sum_{\vec{l}_\beta, \vec{d}_\beta}^{(l)} := \sum_{\beta=1}^l \sum_{\vec{l}_\beta \in L_{l,\beta}} \sum_{\vec{d}_\beta \in \{1, \dots, d\}^\beta} \quad (\text{for } l \geq 1), \quad \sum_{\vec{l}_\beta, \vec{d}_\beta}^{(0)} := \sum_{\beta=0} \sum_{\vec{l}_0 = (\emptyset)} \sum_{\vec{d}_0 = (\emptyset)} \quad (\text{for } l = 0), \quad (3)$$

$$L_{l,\beta} := \left\{ \vec{l}_\beta = (l_1, \dots, l_\beta); \sum_{j=1}^\beta l_j = l, (l, l_j, \beta \in \mathbf{N}) \right\}, \quad (4)$$

$$\sum_{\vec{k}_\delta}^{(n)} = \sum_{\delta=1}^n \sum_{\vec{k}_\delta \in L_{n,\delta}}. \quad (5)$$

Then, the expectation of $\Phi(G^{(\epsilon)})$ is expanded as follows.

Basket Option

- **Basket options are one of popular exotic options in commodity and equity markets.**
 - **The payoff of a basket call option is expressed as $(g(S_T) - K)^+$.**
 - **g is a basket price function defined by $g(S_T) := \sum_{i=1}^d w_i S_T^i$, where $S_T = (S_T^1, \dots, S_T^d)$ are asset prices, and w_i is a constant weight for each i .**
- **However, it is difficult to calculate a basket option price.**
 - **Numerical methods for PDE**
 - **It is difficult to solve high dimensional PDEs.**
 - **Monte Carlo method**
 - **It needs a large amount of computational time to obtain an accurate value.**

Numerical Example: Basket Option (LSV model)

Numerical Example: Basket Option with 100 Underlying Assets (Shiraya-T(2014))

Table: Basket Call Option ($T = 1$)

Strike(K)	8,000	9,000	10,000	11,000	12,000
Monte Carlo	2,037.1	1,167.5	517.6	160.8	31.7
AE3rd	2,037.4	1,167.6	517.6	160.5	31.5
Difference	0.3	0.2	-0.0	-0.2	-0.2
Relative Difference (%)	0.0%	0.0%	0.0%	-0.2%	-0.7%
MC Std Error	0.7	0.6	0.4	0.2	0.2

Monte Carlo: the number of trials is 3 million with the antithetic variable method.

Approximate density of CVA

We show the density function of approximate CVA by the asymptotic expansion method with Monte Carlo simulation. The maturity of forward contract is τ , T denotes the future time when CVA is evaluated, and Δ denotes the lag of collateral.

- maturity (τ): 5 years, evaluation date (T): 2.5 years.
- strike: 10,000.
- time step size: $\frac{1}{400}$ year.
- the number of trials: 325,000 with antithetic variates.

Procedure:

- 1 implement Monte carlo simulation of the state variables (h^1, h^2, S, v) until T .
- 2 given each realization of the state variables, compute $c\hat{v}a_{AE}(T, u)$.
- 3 integrate $c\hat{v}a_{AE}(T, u)$ numerically with respect to the time parameter u from T to τ , and plot the values and their frequencies after normalization.

Perturbation Technique with Interacting Particle Method

- Its dynamics in terms of the future time u is specified by $(Y_{t,u})_j^i = \partial_{x_j} X_u^i$:

$$\begin{aligned} d(Y_{t,u})_j^i &= \partial_k \gamma_0^i(X_u)(Y_{t,u})_j^k du + \partial_k \gamma_a^i(X_u)(Y_{t,u})_j^k dW_u^a \\ (Y_{t,t})_j^i &= \delta_j^i \end{aligned} \tag{93}$$

where ∂_k denotes the differential with respect to the k -th component of X , and δ_j^i denotes Kronecker delta. Here, i and j run through $\{1, \dots, d\}$ and $\{1, \dots, r\}$ for a . Here, we adopt Einstein notation which assumes the summation of all the paired indexes.

- Then, it is well-known that

$$(\mathcal{D}_t X_u^i)_a = (Y_{t,u} \gamma(x_t))_a^i,$$

where $a \in \{1, \dots, r\}$ is the index of r -dimensional Brownian motion.

Perturbation Technique with Interacting Particle Method

- ϵ -1st order:

$$\begin{aligned} V_t^{(1)} &= \int_t^T \mathbb{E}[f(X_u, V_u^{(0)}, Z_u^{(0)}) | \mathcal{F}_t] du \\ &= \int_t^T \mathbb{E}[f(X_u, v^{(0)}(X_u), z^{(0)}(X_u)) | \mathcal{F}_t] du \end{aligned} \quad (97)$$

- Next, define the new process for ($s > t$):

$$\hat{V}_{ts}^{(1)} = e^{\int_t^s \lambda_u du} V_s^{(1)}, \quad (98)$$

where deterministic positive process λ_t (It can be a positive constant for the simplest case.).

Perturbation Technique with Interacting Particle Method

- Then, its dynamics is given by

$$d\hat{V}_{ts}^{(1)} = \lambda_s \hat{V}_{ts}^{(1)} ds - \lambda_s \hat{f}_{ts}(X_s, v^{(0)}(X_s), z^{(0)}(X_s)) ds + e^{\int_t^s \lambda_u du} Z_s^{(1)} \cdot dW_s,$$

where

$$\hat{f}_{ts}(x, v^{(0)}(x), z^{(0)}(x)) = \frac{1}{\lambda_s} e^{\int_t^s \lambda_u du} f(x, v^{(0)}(x), z^{(0)}(x)).$$

- Since we have $\hat{V}_{tt}^{(1)} = V_t^{(1)}$, one can easily see the following relation holds:

$$V_t^{(1)} = \mathbb{E} \left[\int_t^T e^{-\int_t^u \lambda_s ds} \lambda_u \hat{f}_{tu}(X_u, v^{(0)}(X_u), z^{(0)}(X_u)) du \middle| \mathcal{F}_t \right] \quad (99)$$

- As in credit risk modeling (e.g. Bielecki-Rutkowski (2002)), it is the present value of default payment where the default intensity is λ_s with the default payoff at $s(> t)$ as $\hat{f}_{ts}(X_s, v^{(0)}(X_s), z^{(0)}(X_s))$. Thus, we obtain the following proposition.

Perturbation Technique with Interacting Particle Method

Proposition

The $V_t^{(1)}$ in (97) can be equivalently expressed as

$$V_t^{(1)} = \mathbf{1}_{\{\tau > t\}} \mathbb{E} \left[\mathbf{1}_{\{\tau < T\}} \hat{f}_{t\tau} \left(X_\tau, v^{(0)}(X_\tau), z^{(0)}(X_\tau) \right) \middle| \mathcal{F}_t \right]. \quad (100)$$

Here τ is the interaction time where the interaction is drawn independently from Poisson distribution with an arbitrary deterministic positive intensity process λ_t . \hat{f} is defined as

$$\hat{f}_{ts}(x, v^{(0)}(x), z^{(0)}(x)) = \frac{1}{\lambda_s} e^{\int_t^s \lambda_u du} f(x, v^{(0)}(x), z^{(0)}(x)). \quad (101)$$

Perturbation Technique with Interacting Particle Method

- Now, let us consider the component $Z^{(1)}$. It can be expressed as

$$Z_t^{(1)} = \int_t^T \mathbb{E} \left[\mathcal{D}_t f(X_u, v^{(0)}(X_u), z^{(0)}(X_u)) \middle| \mathcal{F}_t \right] du \quad (102)$$

Firstly, let us observe the dynamics of Malliavin derivative of $V^{(1)}$ follows

$$\begin{aligned} d(\mathcal{D}_t V_s^{(1)}) &= -(\mathcal{D}_t X_s^i) \nabla_i(X_s, v^{(0)}, z^{(0)}) f(X_s, v^{(0)}, z^{(0)}) + (\mathcal{D}_t Z_s^{(1)}) \cdot dW_s; \\ \mathcal{D}_t V_t^{(1)} &= Z_t^{(1)}, \end{aligned} \quad (103)$$

where

$$\nabla_i(x, v^{(0)}, z^{(0)}) \equiv \partial_i + \partial_i v^{(0)}(x) \partial_v + \partial_i z^{(0)}(x) \partial_{z^a}, \quad (104)$$

$$f(x, v^{(0)}, z^{(0)}) \equiv f(x, v^{(0)}(x), z^{(0)}(x)). \quad (105)$$

Perturbation Technique with Interacting Particle Method

- Define, for $(s > t)$,

$$\widehat{\mathcal{D}}_t V_s^{(1)} = e^{\int_t^s \lambda_u du} (\mathcal{D}_t V_s^{(1)}). \tag{106}$$

Then, its dynamics can be written as

$$\begin{aligned} d(\widehat{\mathcal{D}}_t V_s^{(1)}) &= \lambda_s (\widehat{\mathcal{D}}_t V_s^{(1)}) ds - \lambda_s (\mathcal{D}_t X_s^i) \nabla_i (X_s, v^{(0)}, z^{(0)}) \hat{f}_{ts}(X_s, v^{(0)}, z^{(0)}) ds \\ &\quad + e^{\int_t^s \lambda_u du} (\mathcal{D}_t Z_s^{(0)}) \cdot dW_s. \end{aligned} \tag{107}$$

We again have

$$\widehat{\mathcal{D}}_t V_t^{(1)} = Z_t^{(1)}. \tag{108}$$

- Hence,

$$Z_t^{(1)} = \mathbb{E} \left[\int_t^T e^{-\int_t^u \lambda_s ds} \lambda_u (\mathcal{D}_t X_u^i) \nabla_i (X_u, v^{(0)}, z^{(0)}) \hat{f}_{tu}(X_u, v^{(0)}, z^{(0)}) du \middle| \mathcal{F}_t \right]. \tag{109}$$

Perturbation Technique with Interacting Particle Method

Monte Carlo Method

Now, we have a new particle interpretation of $(V^{(1)}, Z^{(1)})$ as follows:

$$V_t^{(1)} = \mathbf{1}_{\{\tau > t\}} \mathbb{E} \left[\mathbf{1}_{\{\tau < T\}} \hat{f}_{t\tau}(X_\tau, v^{(0)}, z^{(0)}) \middle| \mathcal{F}_t \right] \quad (111)$$

$$Z_t^{(1)} = \mathbf{1}_{\{\tau > t\}} \mathbb{E} \left[\mathbf{1}_{\{\tau < T\}} (Y_{t,\tau} \gamma(X_\tau))^i \nabla_i (X_\tau, v^{(0)}, z^{(0)}) \hat{f}_{t\tau}(X_\tau, v^{(0)}, z^{(0)}) \middle| \mathcal{F}_t \right] \quad (112)$$

which allows efficient time integration with the following Monte Carlo scheme:

- Run the diffusion processes of X and Y
- Carry out Poisson draw with probability $\lambda_s \Delta s$ at each time s and if "one" is drawn, set that time as τ .
- Then stores the relevant quantities at τ , or in the case of $(\tau > T)$ stores 0.
- Repeat the above procedures and take their expectation.

