

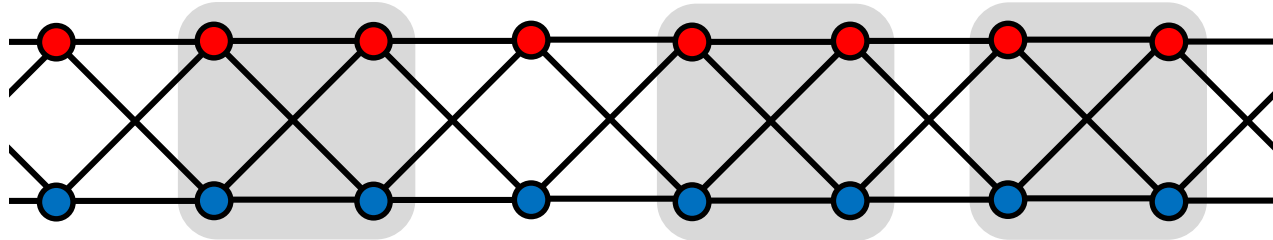
# 平坦バンドのあるBose-Hubbard模型におけるペア朝永-Luttinger液体相

桂 法称 (東大院理)

高吉慎太郎 (東大院理 → Univ. Geneva)

渡辺伯陽 (カブリIPMU)

青木秀夫 (東大院理)



- S. Takayoshi, H.K., N. Watanabe, and H. Aoki, *Phys. Rev. A* **88**, 063613 (2013).

# Creutz ladder

## ■ Some of my favorite subjects:

1. Bethe ansatz
2. Flat bands
3. Topological insulators

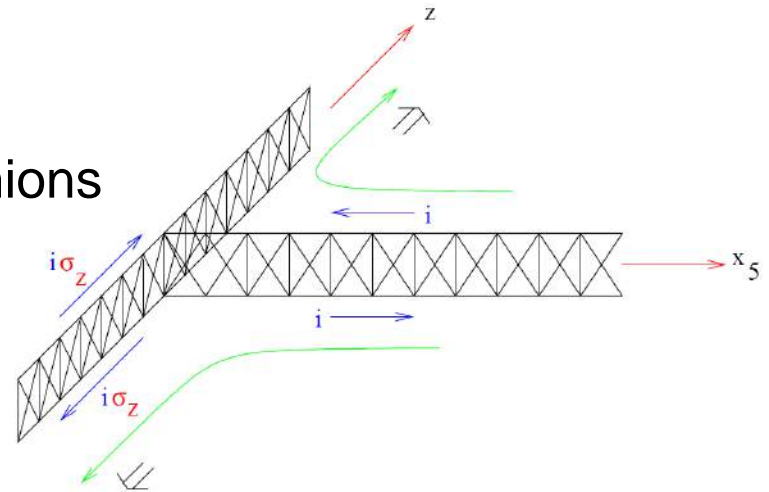
All of them lead to *non-perturbative* results.

Marriage of these 3 in the Creutz ladder → *Pair TL liquids!*

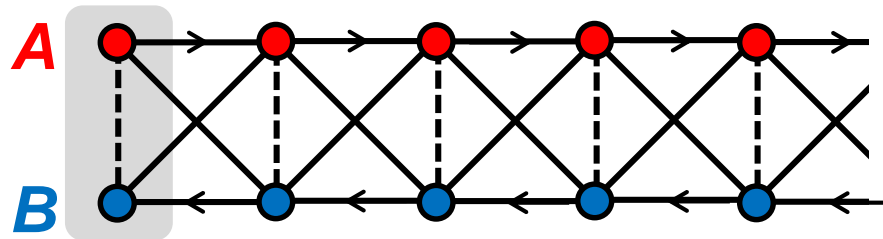
## ■ Creutz ladder (free-fermion model)

M. Creutz, *PRL* **83**, 2636 (1999).

- 1d toy model for domain-wall fermions
- Edge states
- Spin separator?



# What about Bose-Hubbard on the Creutz ladder?



Bosons  $a_j$  live on chain  $A$ , while bosons  $b_j$  on chain  $B$ .

■ Hamiltonian:  $\mathcal{H} = \mathcal{H}_{\text{hop}} + \mathcal{H}_U$

Hopping term:

$$\mathcal{H}_{\text{hop}} = -t \sum_j (e^{i\theta} a_j^\dagger a_{j+1} + e^{-i\theta} b_j^\dagger b_{j+1} + \text{h.c.})$$

$$- t_\times \sum_j (a_j^\dagger b_{j+1} + b_j^\dagger a_{j+1} + \text{h.c.}) - t_\perp \sum_j (a_j^\dagger b_j + \text{h.c.})$$

Class AIII & BDI **topological insulator** (Mazza *et al*, *NJP* **14** ('12))  
Edge states for  $\theta = \pi/2$  and small  $t_\perp$ .

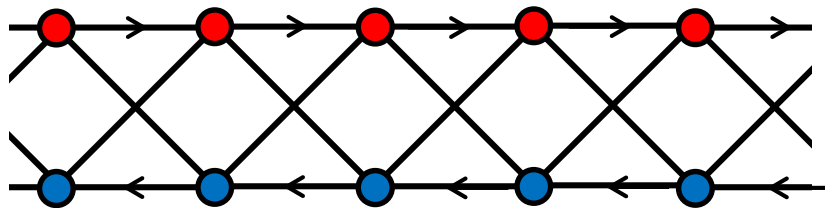
On-site repulsion ( $U > 0$ ):

$$\mathcal{H}_U = \frac{U}{2} \sum_j a_j^\dagger a_j^\dagger a_j a_j + \frac{U}{2} \sum_j b_j^\dagger b_j^\dagger b_j b_j$$

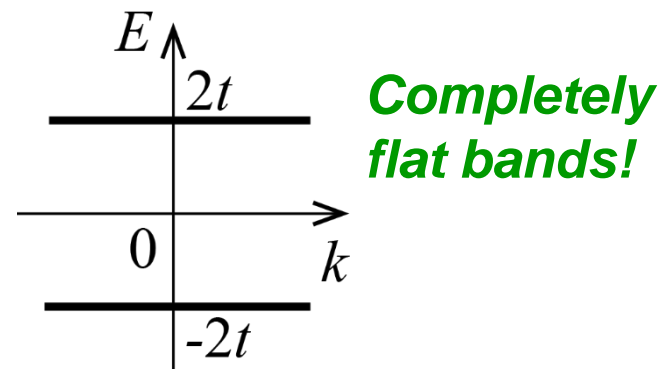
How the interaction affects 1d topo. ins.?  
Fermions v.s. bosons?

# Flat band and Wannier basis

## ■ Flat-band case



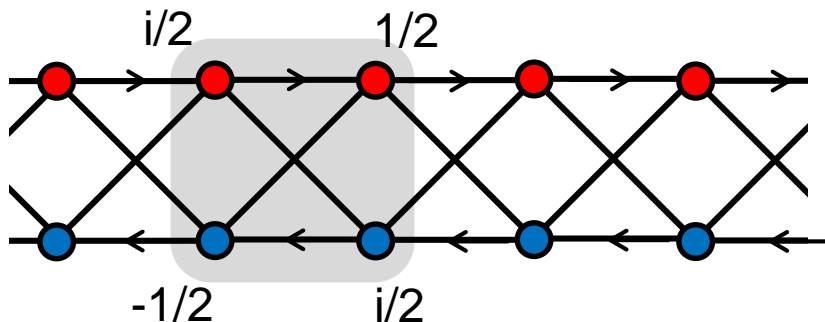
$$t = t_x, \quad t_{\perp} = 0, \quad \theta = \pi/2$$



$$E(k) = \pm 2t$$

## ■ Wannier states

Usually, flat bands are spanned by non-orthogonal basis states. Here, we can construct a **Wannier basis** for both upper and lower flat bands. For the lower band, we have



$$w_j = \frac{1}{2}(\mathbf{i}a_j - b_j - a_{j+1} + \mathbf{i}b_{j+1})$$

$$[w_i, w_j^\dagger] = \delta_{i,j} \quad \text{Orthonormal!}$$

# Wigner solid

- Boson density

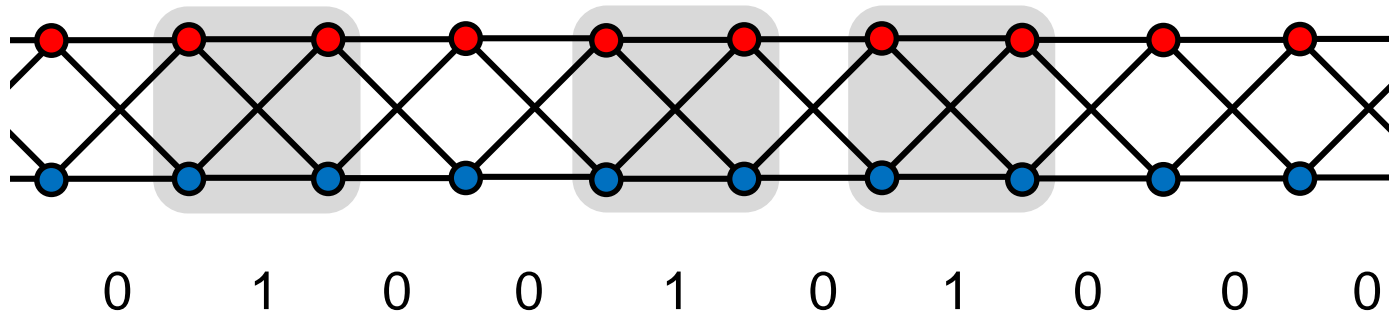
$$\rho := \frac{N}{2L}$$

$L$ : Number of sites on one chain  
 $N$ : Total number of bosons

- Ground states for low densities ( $\rho \leq 1/4$ )

Since  $H_U$  is **positive semi-definite**, any zero-energy state of  $H_U$  spanned by  $w^\dagger$  is an exact ground state of  $H = H_{\text{hop}} + H_U$ .

→ *1-to-1 correspondence with non-overlapping configurations!*



They do not touch each other as long as  $N \leq L/2$ .

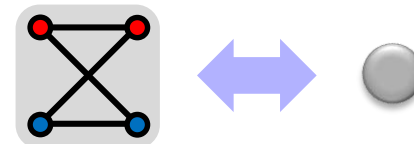
→ **Wigner solid** phase for  $\rho \leq 1/4$ .

# Low-energy effective Hamiltonian ( $U < t$ )

## ■ Projected interaction

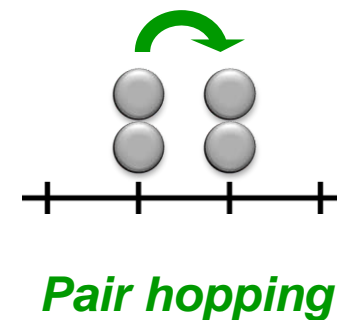
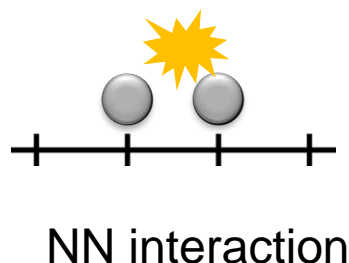
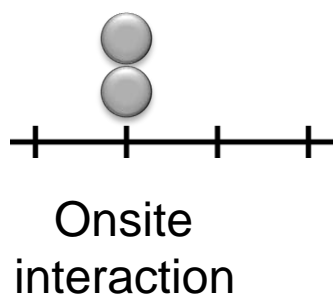
Consider the case  $U \ll 4t$  (band gap).

$H_U$  projected onto the lower band reads



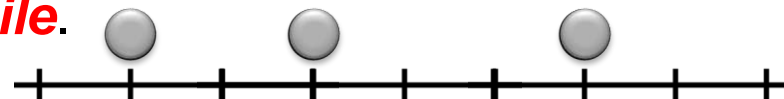
$$\mathcal{H}_{\text{proj}}^{\text{int}} = \sum_{j=1}^L \left[ \frac{U}{8} w_j^\dagger w_j^\dagger w_j w_j + \frac{U}{4} w_j^\dagger w_{j+1}^\dagger w_j w_{j+1} - \frac{U}{16} (w_j^\dagger w_j^\dagger w_{j+1} w_{j+1} + \text{H.c.}) \right]$$

Positive semi-definite



Isolated single bosons are *immobile*.

Wigner solids correspond to

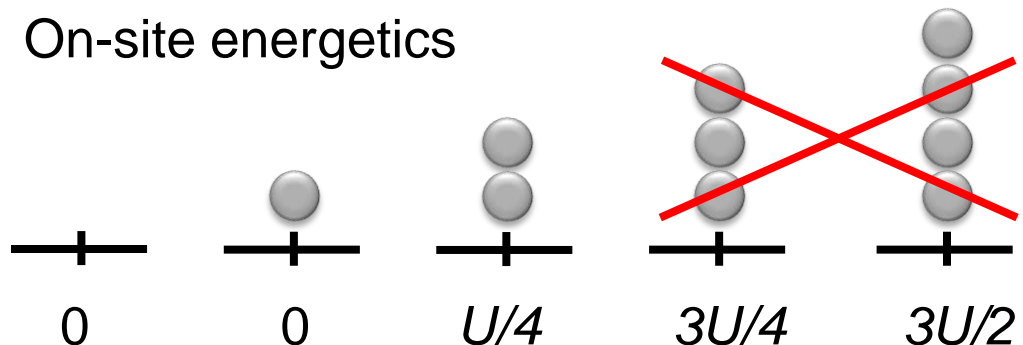


*Pairs of bosons* can lower the “kinetic” energy.

(Similar to 2-magnon bound states in *nematic phases*.)

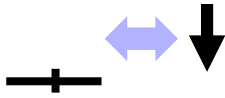
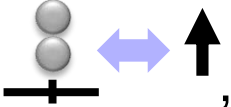
# Effective spin model for high densities

## On-site energetics



- Neglect states with more than 2 bosons at one site.
- On-site int < NN int.

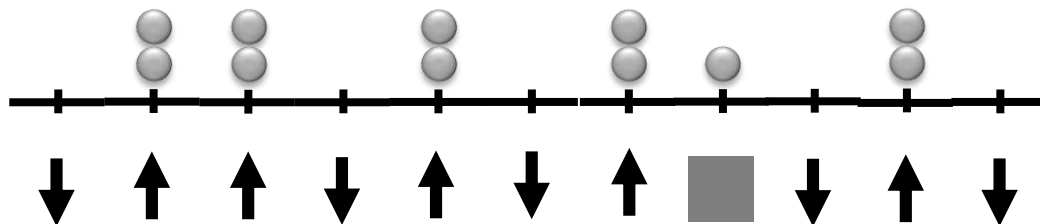
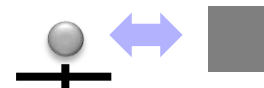
## Spin Hamiltonian

With the identification  and ,

effective spin model = Spin-1/2 XXZ chain in a magnetic field

$$\mathcal{H}_{\text{XXZ}} = \frac{U}{4} \sum_j (-S_j^x S_{j+1}^x - S_j^y S_{j+1}^y + 4S_j^z S_{j+1}^z) + \frac{5}{4}U \sum_j S_j^z$$

Single-boson sites can be thought of as “walls”.



Each configuration is an ensemble of open XXZ chains separated by walls.

# Pair Tomonaga-Luttinger (TL) liquid phase

## ■ Dictionary

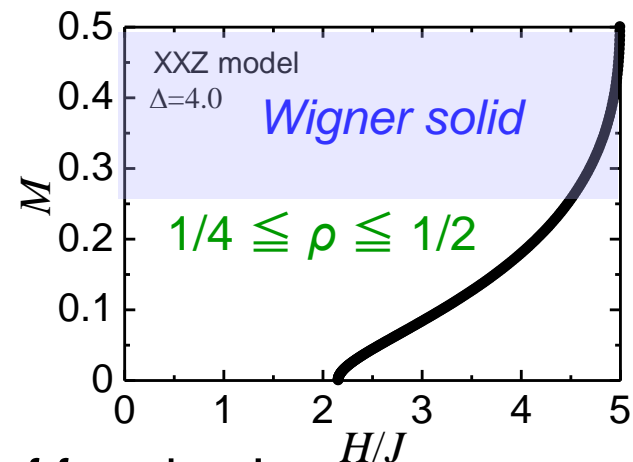
Boson-pair number  $\Leftrightarrow 2S_j^z + 1$

Boson density  $\rho = \frac{1}{2} - M$   $\Leftrightarrow$  Magnetization  $M = -\frac{1}{L} \sum_j S_j^z$

Pair-TL liquid  $\Leftrightarrow$  Field-induced TL liquid

Pairs of bosons are fundamental degrees of freedom!

NOTE: Nematic TL liquid, Doucot-Vidal, *PRL* **88** 227005 ('02)



## ■ Bosonization analysis

TL parameter ( $K$ ) and velocity can be computed via *Bethe ansatz!*

S. Qin *et al.*, *PRB* **56**, 9766 ('97), Cabra *et al.*, *PRB* **58**, 6241 ('98).

- Density-density correlation ( $n_r := w_r^\dagger w_r$ )  $\langle S_r^z S_0^z \rangle$

$$\langle n_r n_0 \rangle \simeq 4\rho^2 - 2K(\pi r)^{-2} + c_1 \cos(2\pi\rho r) r^{-2K}$$

$$1/4 \leq K \leq 1/2$$

- Off-diagonal correlation  $\langle S_r^+ S_0^- \rangle$

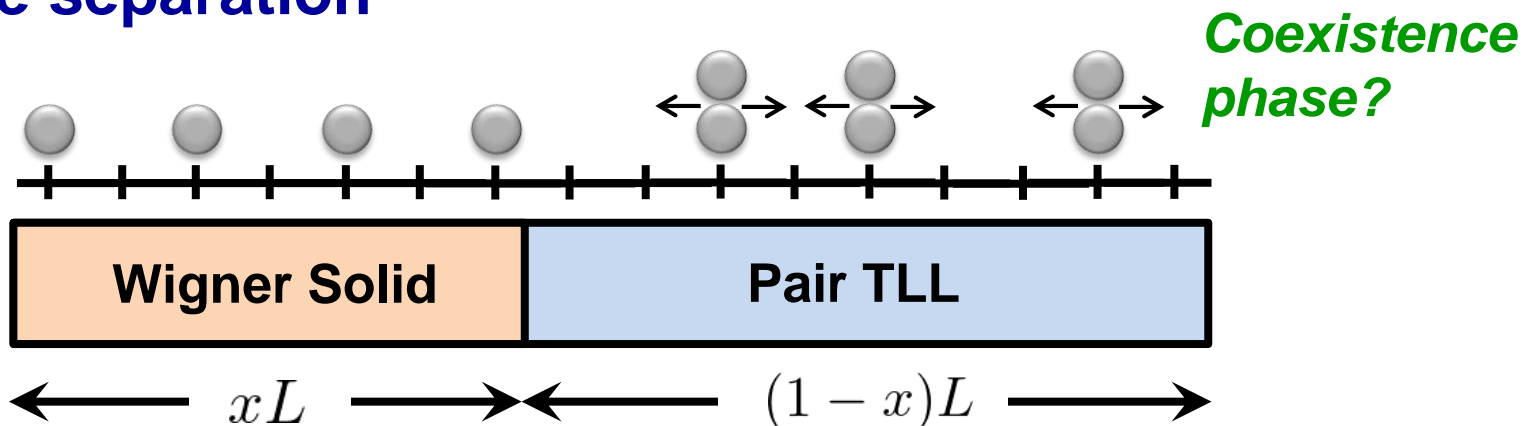
$$\langle w_r^\dagger w_r^\dagger w_0 w_0 \rangle \simeq (-)^r c_2 r^{-\frac{1}{2K}} + (-)^r c_3 \cos(2\pi\rho r) r^{-2K - \frac{1}{2K}}$$

→ The latter decays faster.

→ No supersolid?



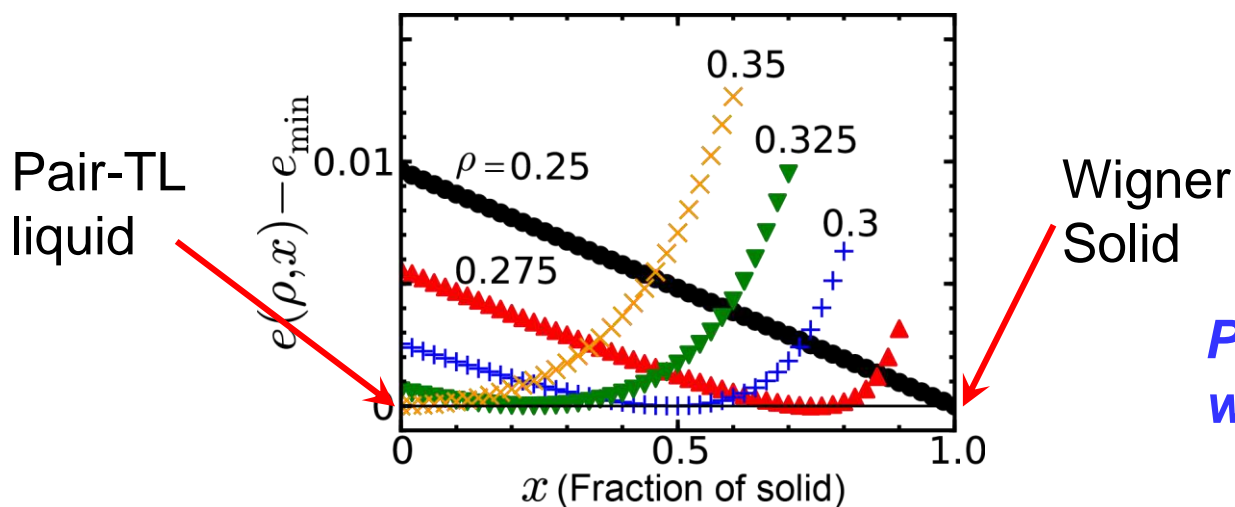
# Phase separation



Density:  $\rho_{\text{WS}} = 1/2$        $\rho_{\text{TLL}} = \frac{\rho - x/4}{1 - x}$

Energy:  $0$        $(1 - x)L e_{\text{TLL}}(M(x))$

Can be computed by solving Bethe-ansatz integral eq.



**Phase separation when  $0.25 < \rho \leq 0.35$ .**

# Charge-density wave at $\rho=1/2$

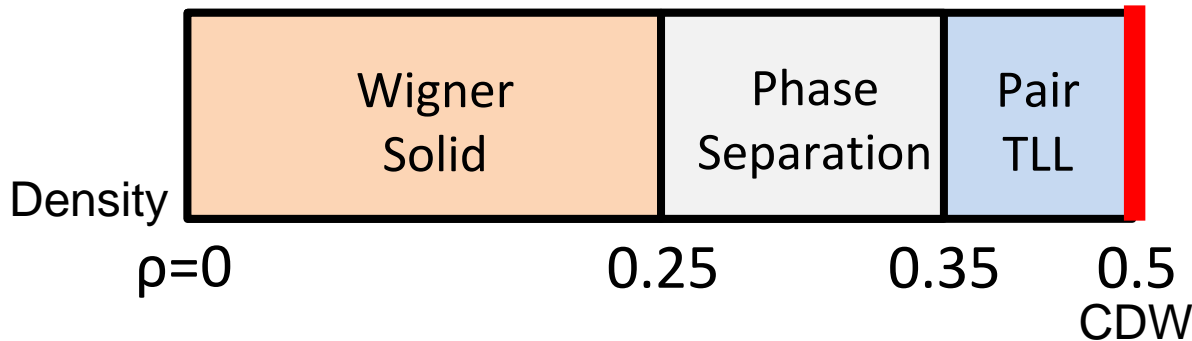
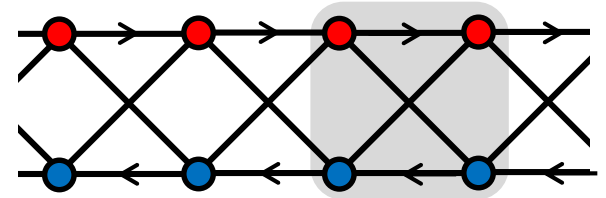
■ Ising-like XXZ  $\rightarrow$  Neel ground states



The 1st and 2nd excited energies can be computed **exactly!**  
 Batchelor-Hamer, *JPA* **23** ('90), Kapustin-Skorik, *JPA* **29** ('96).

## Summary

- Studied Bose-Hubbard on Creutz ladder
- Flat bands and **Wigner solid** ( $\rho \leq 1/4$ )
- Phase diagram for  $U \ll 4t$  (band gap)  $\rightarrow$  **Pair-TL liquid!**



The validity of the XXZ description was checked (ED).

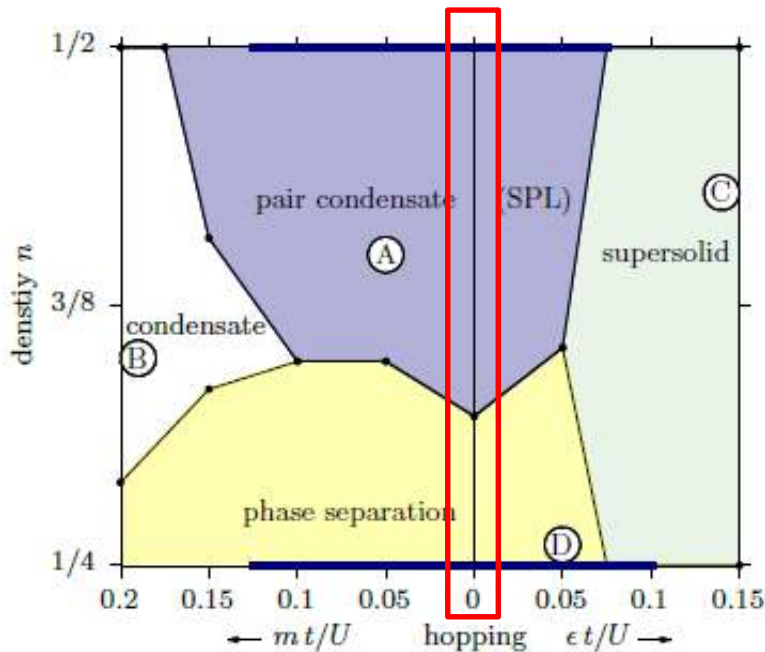
# Away from the flat-band limit

- DMRG study (after our work)

Tovmasyan, Nieuwenburg & Huber, *PRB* **88**, 220510(R) (2013).

Effective Hamiltonian

$$\mathcal{H}_{\text{eff}} = \sum_{j=1}^L \left[ \frac{U}{8} w_j^\dagger w_j^\dagger w_j w_j + \frac{U}{4} w_j^\dagger w_{j+1}^\dagger w_j w_{j+1} - \frac{U}{16} (w_j^\dagger w_j^\dagger w_{j+1} w_{j+1} + \text{H.c.}) \right. \\ \left. + \frac{mt}{4} w_j^\dagger w_{j+1} + \frac{\epsilon t}{4} w_j^\dagger w_{j+2} + \text{H.c.} \right] \quad (mt = t_\perp, \quad \epsilon t = t_\times - t)$$



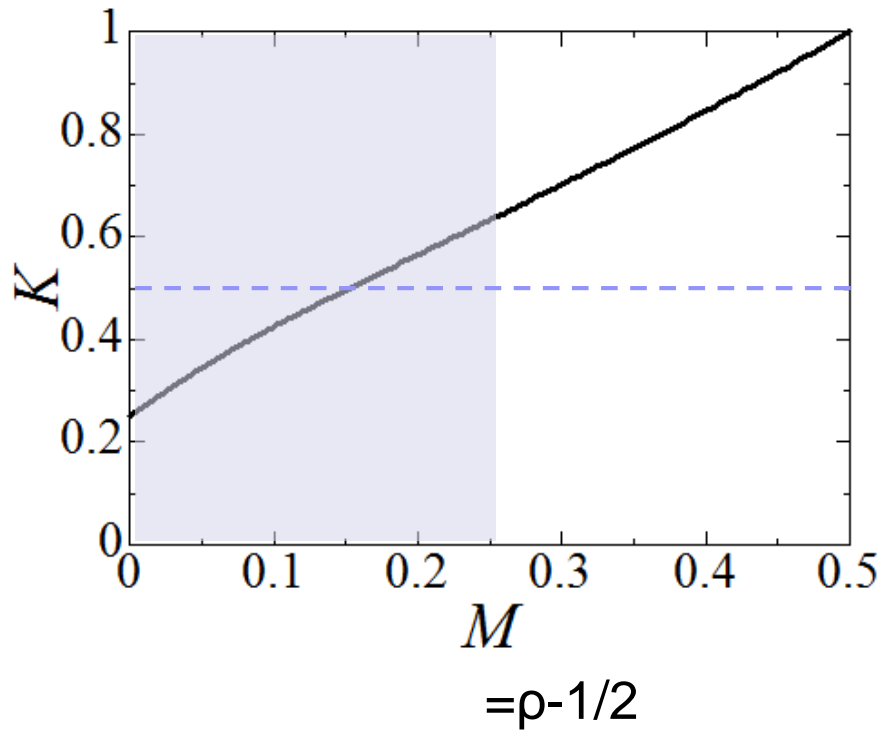
**Qualitatively similar result!**  
**Pair-TL liquid phase is stable**  
**against small perturbations.**

## Future directions

- Realization using synthetic gauge?  
 Y.-J. Lin *et al.*, *Nature* **462**, 628 ('09).
- Fermionic models, topological phases?  
 Spin-orbit interpretation is possible.

# Backup slide

## ■ TL parameter



## ■ TL exponents

