

量子トライマー模型のRVB状態と Z_3 トポロジカル秩序

桂 法称 (東大院理)



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➤ *Phys. Rev. B*, **95**, 060413(R) (2017) [arXiv:1612.06899]

Resonating valence bond (RVB) state

■ What are dimers?

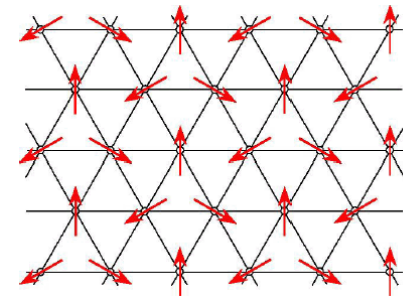
$$\text{dimer } i-j = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Dimer = spin singlet
= valence bond

■ What is RVB?

- $S=1/2$ Heisenberg AFM model on Δ lattice

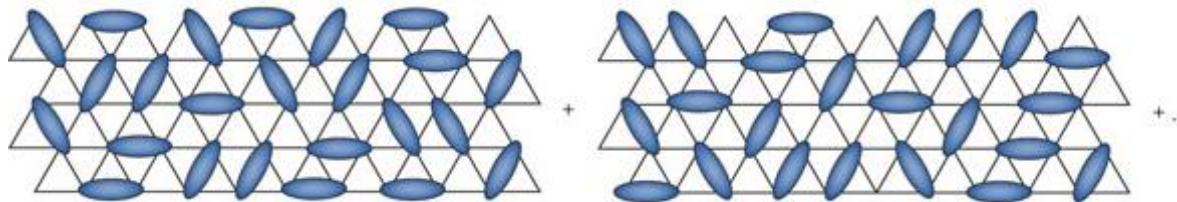
$$H = J \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j$$



Classically, the g.s. exhibits 120° order. What about quantum?

RVB = Equal-weight superposition of all dimer coverings

P.W. Anderson, *Mat. Res. Bull.* **8**, 153 (1973).



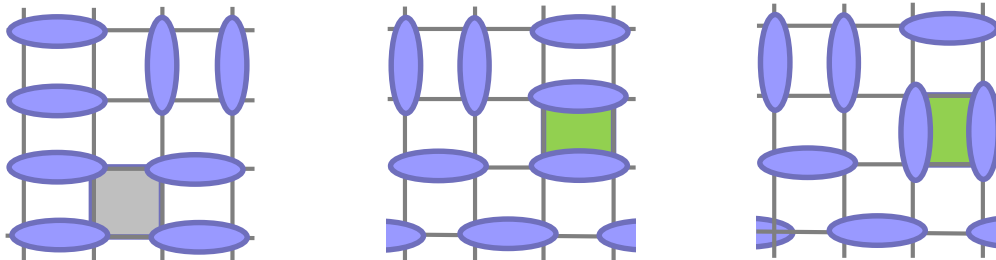
Common belief: This is unlikely. The g.s. has 120° order.

Quantum dimer model

■ Model

Rokhsar-Kivelson, *PRL* **61**, 2376 (1988).

- Basis states: dimer coverings on a square lattice



Dimers don't touch/overlap.
Different configurations are orthogonal.

$$\langle C|C' \rangle = \delta_{C,C'}$$

- Hamiltonian

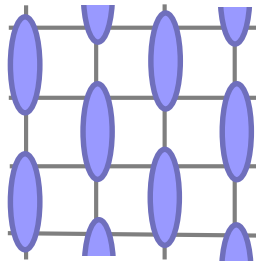
$$\mathcal{H}_{\text{dimer}} = \sum_{\text{plaquettes}} [-J(|\parallel\rangle\langle =| + \text{H.c.}) + V(| =\rangle\langle =| + |\parallel\rangle\langle \parallel|)]$$

kinetic

potential

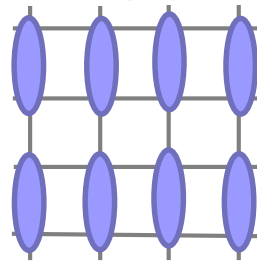
■ Ground state

- $V \gg J > 0$



staggered

- $V < 0 (|V| \gg J)$



columnar

- $V = J$ (RK point) **RVB state!**

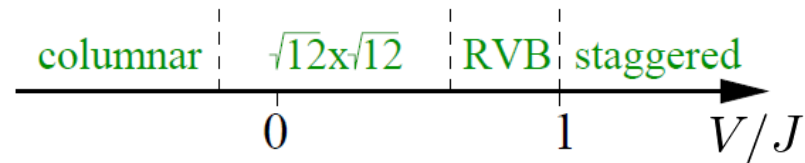
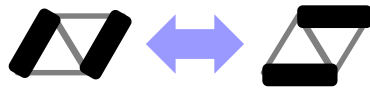
Exact g.s.: $|\Psi_0\rangle = \sum_{C: \text{flippable}} |C\rangle$

Critical dimer-dimer correlation at RK \rightarrow Critical (gapless) point

Topological order

The model on a triangular lattice exhibits *topological order!*

Moessner-Sondhi, *PRL* **86** (2001), Ivanov, *PRB* **70** (2004).

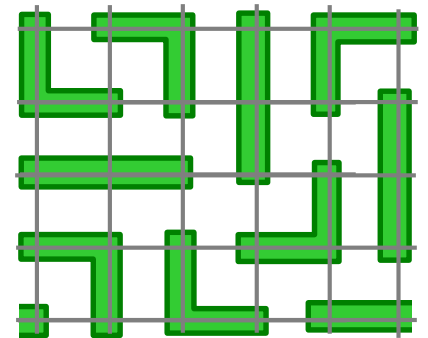


■ What is topological order?

- Spectral gap above the g.s.
gap in RVB phase, vison excitations
- G.S. degeneracy depends on topology
4-fold deg. on a torus
- All g.s. are indistinguishable locally

■ Beyond dimers?

- RVB states consisting of trimers
- RK type Hamiltonian and exact ground states
- Topological degeneracy, \mathbf{Z}_3 vortex excitations



Trimer covering

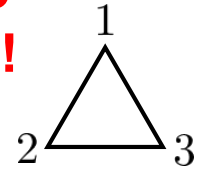
■ What are trimers?

Trimer = $SU(2)$ singlet made up of three $S=1$

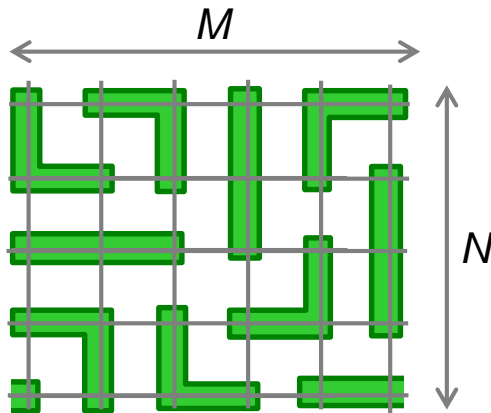
$$1 \otimes 1 \otimes 1 = 3 \oplus 2 \oplus 2 \oplus 1 \oplus 1 \oplus 1 \oplus 0$$

0 is the g.s. of 3-site AFM Heisenberg chain

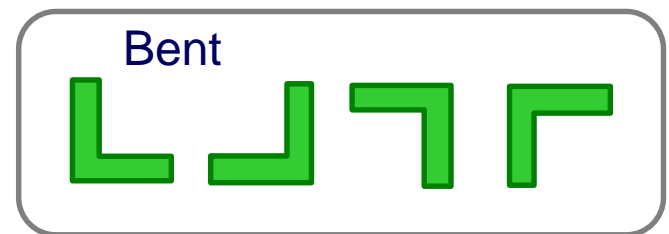
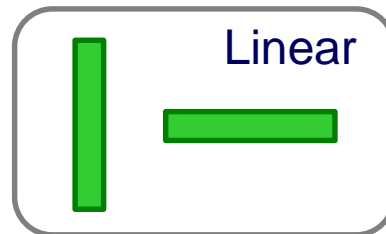
**Unique
singlet!**



■ Trimers on square lattice



• Allowed trimers



• Rules

- Place trimers without making holes
- Trimers should not touch or overlap

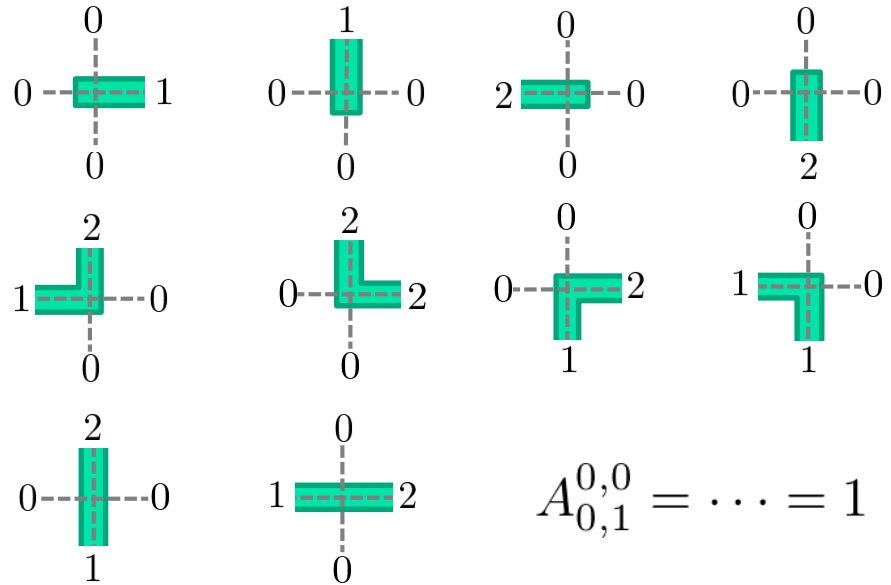
Q. How many ways to arrange trimers?

Tensor network approach

Local tensor

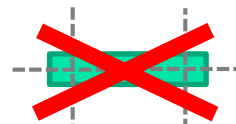
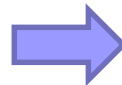
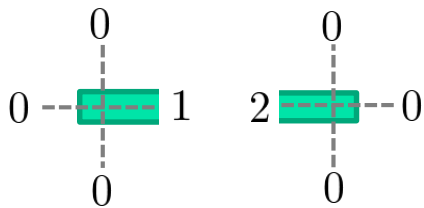
$$A_{\gamma,\delta}^{\alpha,\beta} = \begin{array}{c} \alpha \\ | \\ \gamma \text{---} \text{---} \delta \\ | \\ \beta \end{array}$$

$\alpha, \beta, \gamma, \delta = 0, 1, \text{ or } 2$ labels
a state on each **edge**.
Only 10 nonzero elements.

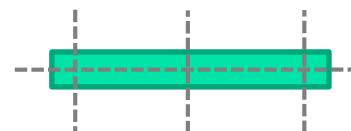
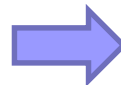
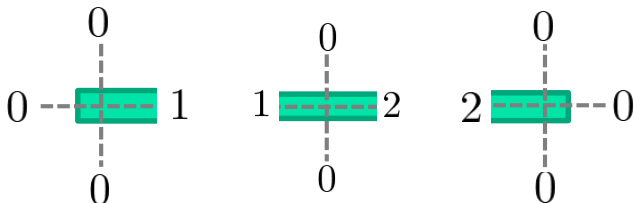


$$A_{0,1}^{0,0} = \dots = 1$$

How it works



Dimers never appear!



Trimers arise naturally.

Results

■ Counting

Local tensor \rightarrow transfer matrix T_M

$$\#(\text{configs.}) = Z = \text{Tr} (T_M)^N$$

Z	$N=1$	2	3	4	5	6
$M=3$	3	33	174	585	2,598	11,550
6	3	297	11,550	54,417	705,708	9,027,000
9	3	2,913	1,094,943	7,111,413	325,897,458	15,280,181,589

Z grows exponentially with system size!

■ Entropy/site

$$s = \frac{1}{MN} \ln Z \sim \frac{1}{M} \ln \lambda_M^{\max}$$

λ_M^{\max} : Largest eigenvalues of T_M

For large M and N , we have

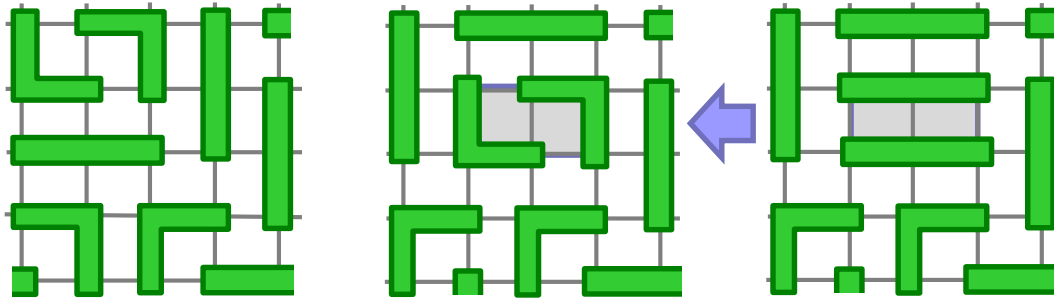
$$s = \begin{cases} 0.41194 & (= \ln 1.50974) \\ 0.27693 & (= \ln 1.31907) \\ 0.15852 & (= \ln 1.17178) \end{cases}$$

bent only Froboese *et al*, *JPA* **29** (1996)
linear only Ghosh *et al*, *PRE* **75** (2007)

Quantum trimer model (1)

■ Basis states

Trimer coverings on a square lattice



Trimers don't touch/overlap.
Different configurations
are orthogonal.

$$\langle \mathcal{T} | \mathcal{T}' \rangle = \delta_{\mathcal{T}, \mathcal{T}'}$$

■ Rokhsar-Kivelson Hamiltonian

$$H = v \left\{ 2 \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| \right.$$

$$\left. + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \dots \right\}$$

Potential term

$$-t \left\{ \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| \right.$$

$$\left. + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| \right.$$

$$\left. + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + \left| \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right\rangle \left\langle \begin{array}{cc} \text{---} & \text{---} \\ \text{---} & \text{---} \end{array} \right| + R_{\frac{\pi}{2}} + h.c. \right\}$$

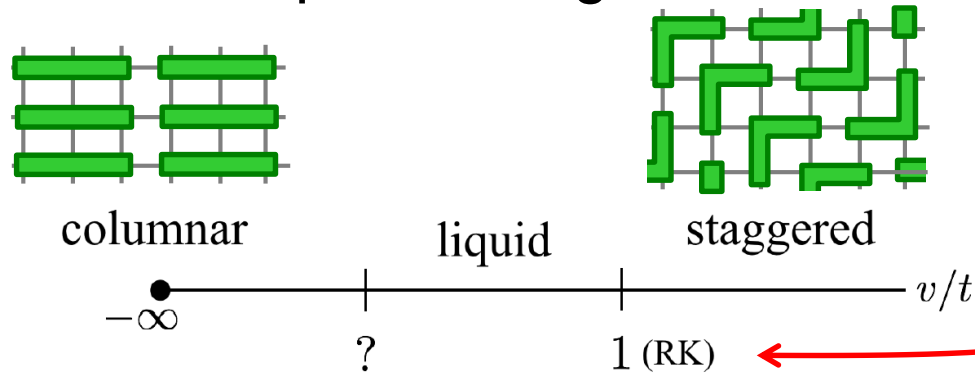
Kinetic (resonance) term

Assume $t > 0$.

Resonance involves
only two trimers.

Quantum trimer model (2)

■ Schematic phase diagram



Exact g.s. at RK point

$$|\text{tRVB}\rangle = \sum_{\mathcal{T}: \text{flippable}} |\mathcal{T}\rangle$$

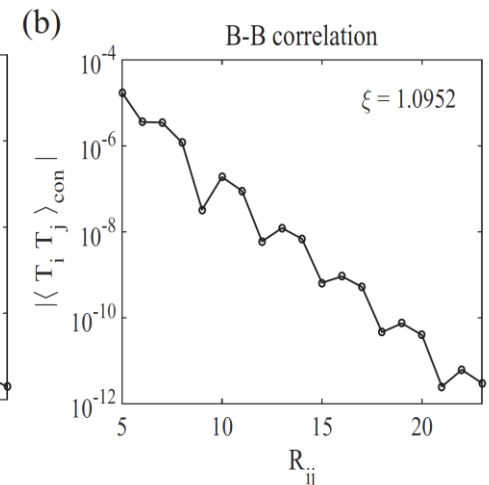
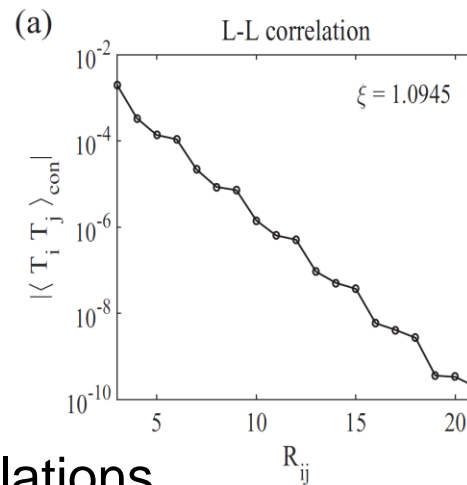
Trimer RVB state!

■ Ground-state correlations in tRVB

$$\langle \text{tRVB} | T_i T_j | \text{tRVB} \rangle \propto \frac{Z'_{ij}}{Z}$$

Z : Total # of trimer configs.

Z'_{ij} : # of configs with fixed trimers at i and j

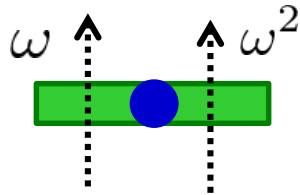


Exponentially decaying correlations
Imply **gapped** nature of the model

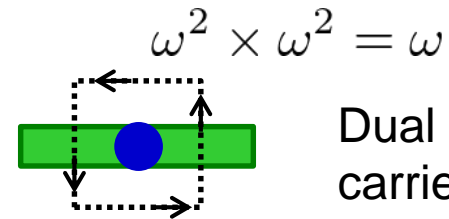
$\xi \sim$ lattice spacing

Topological sectors

Z3 flux



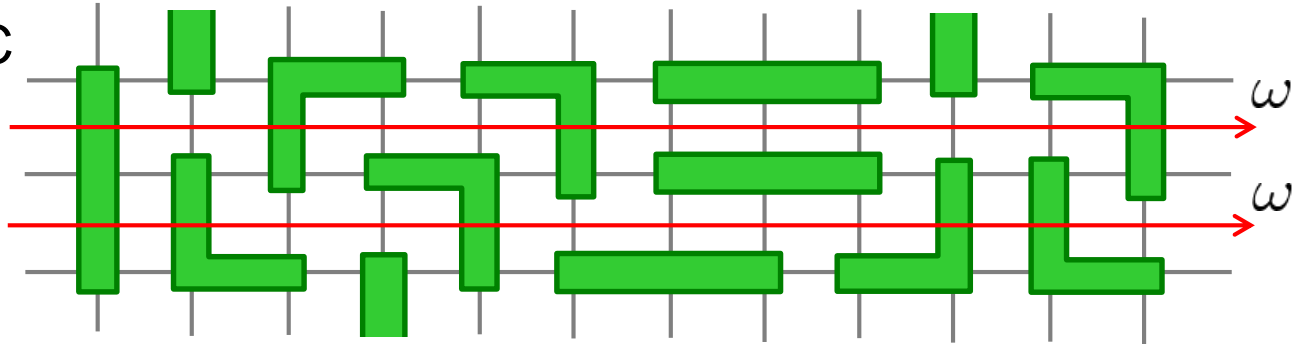
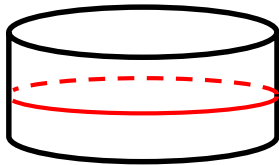
ω when ● is seen on the right. ω^2 when it is seen on the left



Dual plaquette carries flux ω

■ Winding numbers

Cylinder with PBC

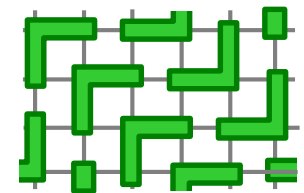


$V_{\Gamma} := \omega^{n_l - n_r}$ commutes with H . $V_{\Gamma} = 1, \omega$ or ω^2 . \rightarrow 3 sectors!

On a torus, we have $3 \times 3 = 9$ disconnected sectors.

■ Ergodicity ...

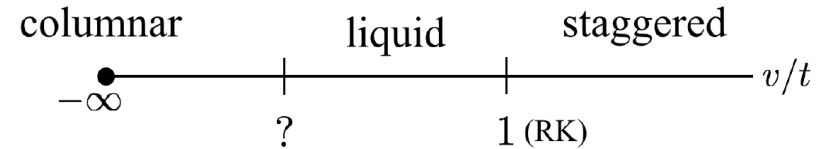
Hamiltonian H is block-diagonal w.r.t. the sectors.
Is the action of H ergodic in each sector? \rightarrow NO!



Staggered states are frozen...

Z3 topological order

■ Around RK point



At RK, $|\text{tRVB}\rangle = \sum_{\mathcal{T}: \text{flippable}} |\mathcal{T}\rangle$ is the exact $E=0$ g.s. in each sector.

Perturb a little bit! $v = t - \epsilon$

Rule out staggered states but still have **9-fold degeneracy**.

Clear sign of **topological order!** (NOTE: unique g.s. with OBC)

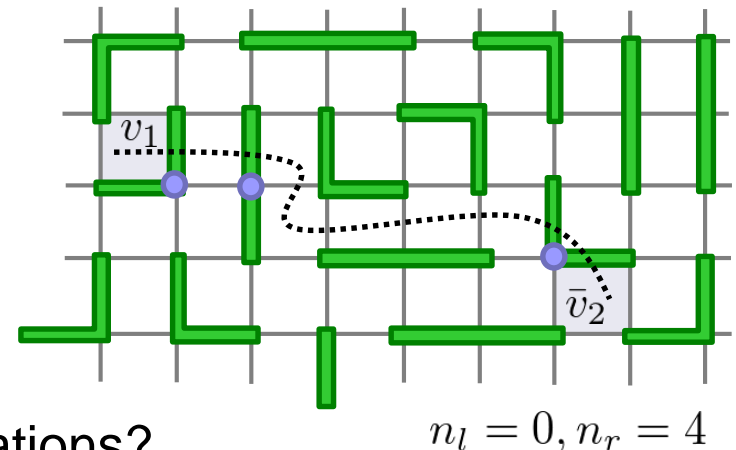
Higher genus surfaces $\rightarrow 9^g$ -fold deg.

■ Z3 vortex excitations

- Variational state

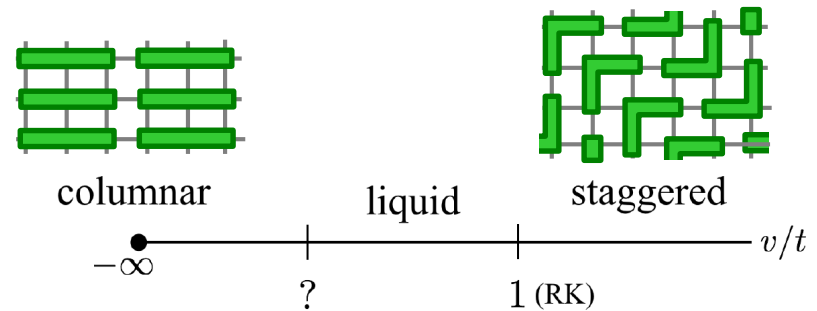
$$|v_1, \bar{v}_2\rangle = \sum_{\mathcal{T}} \omega^{n_l(\mathcal{T}) - n_r(\mathcal{T})} |\mathcal{T}\rangle$$

- Orthogonal to the g.s. $|\text{tRVB}\rangle$
- Close to the true excited states?
- Can the pair split into fractional excitations?

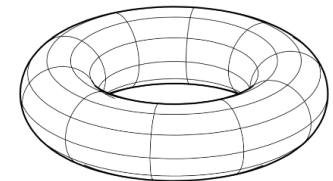


Summary

- Trimer covering on square lattice (entropy/site $s \sim 0.41194$)
- Quantum trimers on square lattice
- Exact ground states at RK point
- Short-range correlation in tRVB
- Topological order

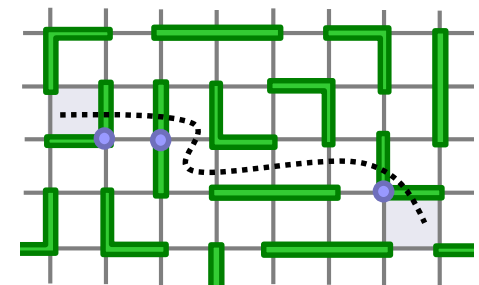


- ❑ Spectral gap above the g.s.
- ❑ 9-fold degeneracy on a torus
- ❑ Z3 vortex excitations?

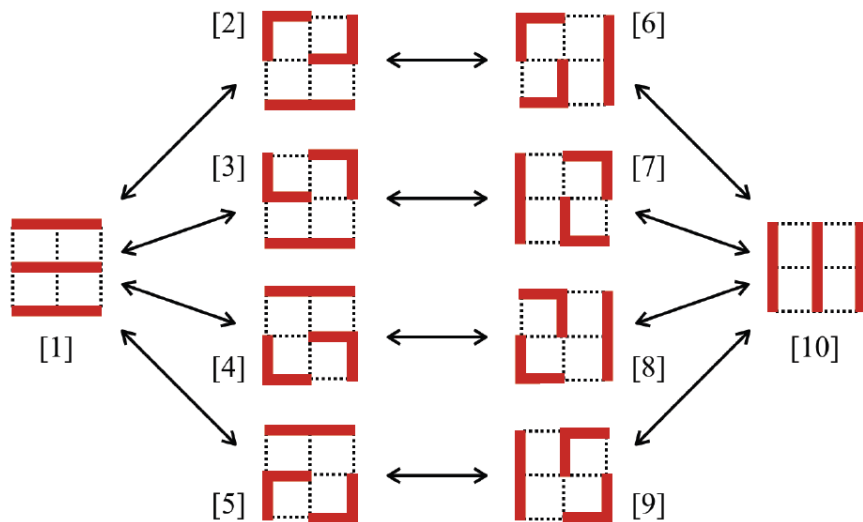


Future directions

- Rigorous proof of exponential decay in tRVB
- Work out phase diagram (ED, QMC, TN, ...)
- Precise nature of the gapped excitations
- Connection to Z3 toric code? Topological QC?



Model with OBC

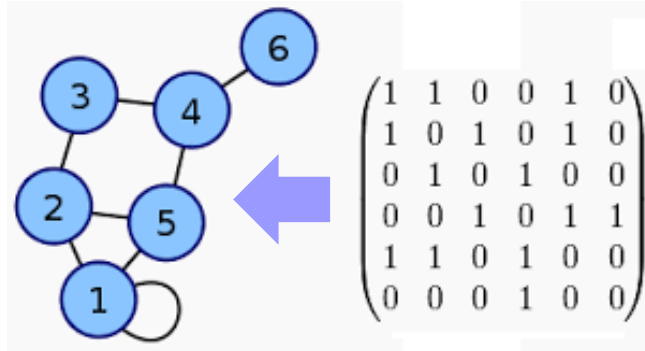


$$H = \begin{bmatrix} 4 & -1 & -1 & -1 & -1 & 0 & 0 & 0 & 0 & 0 \\ -1 & 2 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 2 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ -1 & 0 & 0 & 2 & 0 & 0 & 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & -1 & 0 \\ 0 & -1 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 & 0 & 0 & 2 & 0 & 0 & -1 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 2 & 0 & -1 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 2 & -1 \\ 0 & 0 & 0 & 0 & 0 & -1 & -1 & -1 & -1 & 4 \end{bmatrix}$$

Unique $E=0$ g.s.
at least in this example

Graph analysis

Adjacency matrix/graph



- Graph \rightarrow Matrix

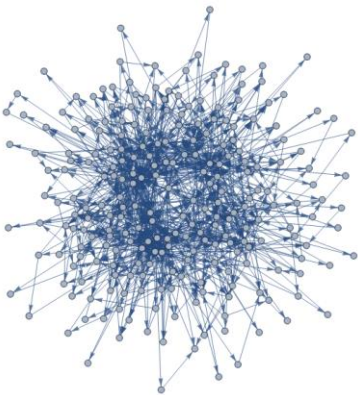
$$A = (a_{ij}), \quad a_{ij} = \begin{cases} 1 & i \sim j \\ 0 & \text{otherwise} \end{cases}$$

- Matrix \rightarrow Graph

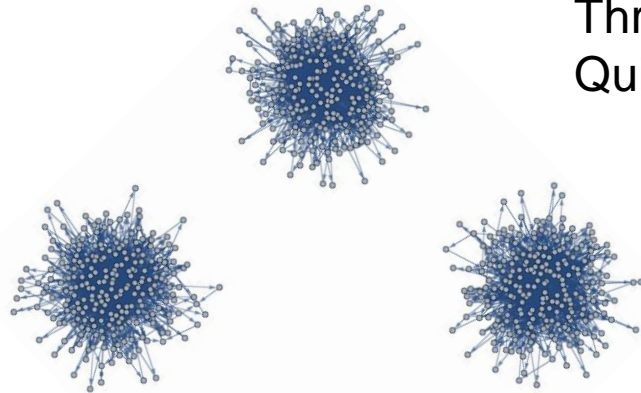
Connect vertices i and j if $a_{ij} \neq 0$.

Adjacency graph of transfer matrix

$M=5$ (3^5 states)



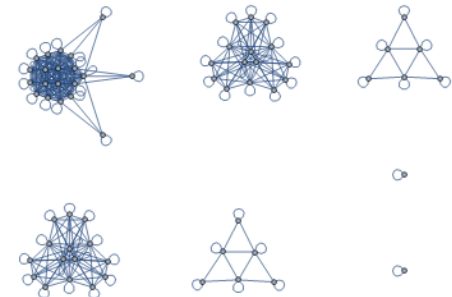
$M=6$ (3^6 states)



Not sparse. Unlikely to be solvable...

Single cluster for $M=3k+1$ or $3k+2$.
Three disconnected clusters for $M=3k$.
Quite different from the dimer case.

Dimer case: $M=6$



General $M \rightarrow M+1$ clusters.